

7N-75 198978 P-20

TECHNICAL NOTE

TEMPERATURE AND COMPOSITION OF A PLASMA OBTAINED BY

SEEDING A CYANOGEN-OXYGEN FLAME WITH CESIUM

By Richard A. Hord and J. Byron Pennington

Langley Research Center Langley Field, Va.

NATIONAL AERONAUTICS AND SPACE ADMINISTRATION
WASHINGTON
May 1960

(NASA-TH-D-38C) TEMPERATURE AND COMPOSITION OF A PLASMA OFFAIRED BY STEDING A CYANGGEN-OXYGEN FLAME WITH CESIUM (NASA-Langley Research Center) 20 p

N89-70877

Unclas 00/75 0198975

NATIONAL AERONAUTICS AND SPACE ADMINISTRATION

TECHNICAL NOTE D-380

TEMPERATURE AND COMPOSITION OF A PLASMA OBTAINED BY

SEEDING A CYANOGEN-OXYGEN FLAME WITH CESIUM

By Richard A. Hord and J. Byron Pennington

SUMMARY

The temperature and composition of a cyanogen-oxygen flame seeded with cesium are calculated approximately by assuming that the combustion is adiabatic and that the product mixture is in thermodynamic equilibrium. For initial mole fractions of cesium up to 0.1 and flame pressures from 0.01 to 1 atmosphere, electron concentrations up to 4.4×10^{16} per cubic centimeter and plasma temperatures in the range of $4,000^{\circ}$ to $5,000^{\circ}$ K are obtained.

INTRODUCTION

A plasma with a sufficiently high concentration of electrons for experimental work is of interest in connection with magneto-gas-dynamic research. Probably the best chemical means available for producing a plasma consists of seeding a high-temperature flame with a gas whose ionization potential is low. It was suggested by Robert V. Hess of the Gas Dynamics Branch of Langley Research Center that seeding the combustion products of cyanogen with alkali metals should yield very high electron concentrations. Subsequent unpublished, approximate calculations performed by Hess and Stuart Patterson showed that electron concentrations of the order of 10^{16} per cubic centimeter could be obtained even with potassium as the seeding material.

It is shown in reference 1 that temperatures in the neighborhood of $5,000^{\circ}$ K are produced in the combustion of a stoichiometric mixture of cyanogen and oxygen. The flame has nearly the highest temperature known to be obtainable by a chemical reaction. The postcombustion mixture consists chiefly of carbon monoxide CO and nitrogen N_2 . Since the physical properties of the CO molecule are generally similar to those of the N_2 molecule, the mixture closely resembles pure nitrogen and, therefore, is to some extent similar to air. A convenient choice of material with low ionization potential is the alkali metal cesium.

The present paper extends the nonseeded, cyanogen-oxygen flame-temperature calculations of reference 1 to cases for which the total pressure is 0.01 to 1 atm and for which the initial mole fraction of cesium is from 0 to 0.1. The partial pressures of the principal species present in the postcombustion mixture are presented along with the electron concentration.

SYMBOLS

е	base of natural logarithms			
μI	precombustion enthalpy per unit mass of mixture			
hII	postcombustion enthalpy per unit mass of mixture			
Hi	enthalpy per mole of particles of the ith species			
h	Planck's constant, 6.62517×10^{-27} erg sec			
I	ionization energy per mole of Cs			
I'	ionization energy per atom of Cs			
k	Boltzmann's constant, $1.38044 \times 10^{-16} \text{ erg/}^{\circ}\text{K}$			
K _i	equilibrium constant (pressure) for dissociation ($i = 1, 2, 3$) or ionization ($i = 5$) of the ith species			
^m i	mass of molecules of ith species			
$^{\rm M}$ i	mass per mole of ith species			
\mathbf{I}^{M}	mass per mole of precombustion mixture			
$^{ m M}$ II	mass per mole of postcombustion mixture			
n _i	number of moles of ith species per unit mass of mixture			
N_{e}	number of electrons per unit volume			
p _i	partial pressure of ith species			
p	total pressure			

fictitious pressure $p_{\mathbf{f}}$ partition function for one particle of ith species q_i (elec) electronic partition function for the cesium atom R universal gas constant per mole \mathbf{T} absolute temperature boiling point of cesium at total pressure p $T_{\mathbf{h}}$ Ui internal energy per mole of ith species U₅(elec) internal energy per mole of cesium due to electronic excitation ٧ volume initial (precombustion) mole fraction of Cs α excitation energy of the jth electronic state of the cesium €.j atom number of particles of ith species in volume $v_{\mathtt{i}}$ density of postcombustion mixture Subscripts: C₂N₂, cyanogen λ 02, diatomic oxygen μ CN, cyanogen radical 1 No, diatomic nitrogen 2 3 CO, carbon monoxide 4 C, atomic carbon

5

Cs, atomic cesium

- 6 N, atomic nitrogen
- 7 0, atomic oxygen
- 8 e⁻, free electron
- 9 Cs⁺, cesium ion

ASSUMPTIONS

The calculation contained herein is based upon the following assumptions:

Cyanogen $\rm C_2N_2$ and oxygen $\rm O_2$ are introduced at room temperature (298.16° K) in stoichiometric proportion corresponding to the highly exothermic chemical reaction:

$$C_2N_2 + O_2 \rightarrow 2CO + N_2$$

Cesium is introduced in monatomic vapor form at a temperature T_b , the boiling point of liquid cesium corresponding to the flame total pressure. The low dissociation energy (0.49 electron volt) of Cs₂ together with the temperatures (e.g., $T_b = 986^{\circ}$ K at p = 1 atm) and the pressures concerned assure the preponderance of the monatomic form in the cesium vapor introduced.

It is also assumed that: (1) the gases are thermally perfect, (2) the enthalpy per unit mass of mixture is the same after combustion as before combustion, (3) the gas mixture reaches thermodynamic equilibrium, and (4) the total pressure is the same after combustion as before combustion.

INITIAL ENTHALPY OF MIXTURE

The components of the precombustion mixture are listed in table I.

TABLE I

Molecule	Mole fraction introduced	Temperature	
C ₂ N ₂	$\frac{1}{2}(1 - \alpha)$	298.16 ^о к	
02	$\frac{1}{2}(1 - \alpha)$	298 . 16 ⁰ к	
Cs	α	$\mathtt{T}_{\mathtt{b}}$	

Hence, the mass per mole of precombustion mixture is

$$M_{\rm I} = \frac{1}{2}(1 - \alpha)M_{\lambda} + \frac{1}{2}(1 - \alpha)M_{\mu} + \alpha M_{5}$$
 (1)

The number of moles of cyanogen ${\rm C_2N_2}$ introduced per unit mass of mixture is then

$$n_{\lambda} = \frac{\frac{1}{2}(1 - \alpha)}{M_{I}}$$

Consequently, the enthalpy (per unit mass of mixture) of cyanogen introduced is

$$n_{\lambda} H_{\lambda} = \frac{\frac{1}{2}(1 - \alpha) H_{\lambda}}{M_{\tau}}$$

where H_{λ} is the enthalpy per mole of C_2N_2 at 298.16° K. The enthalpies (per unit mass of mixture) of oxygen O_2 and cesium Cs can be derived in a similar way.

Therefore, the total enthalpy per unit mass of precombustion mixture is

$$h_{I} = \frac{1}{M_{T}} \left[\frac{1}{2} (1 - \alpha) H_{\lambda} + \frac{1}{2} (1 - \alpha) H_{\mu} + \alpha H_{5} \right]$$
 (2)

where ${\rm H}_{\lambda}$ and ${\rm H}_{\mu}$ correspond to 298.16° K and ${\rm H}_{5}$ corresponds to the temperature ${\rm T}_{b}.$

It is convenient to refer the enthalpy of cesium to the ideal gas enthalpies of the Cs^+ ions and the free electrons at 0° K. Then, the enthalpy per mole of Cs at temperature $\,T_b\,$ is adequately approximated by setting

$$H_5 = \frac{5}{2} RT_b - I$$
 (3)

where I is the ionization energy per mole of Cs. (The corresponding enthalpies, per mole, of Cs⁺ and e⁻ are then $\frac{5}{2}$ RT each, where T is the temperature of the postcombustion mixture.) The molar enthalpies of C_2N_2 and O_2 must be consistent in reference levels with the enthalpies of the postcombustion species CN, N_2 , CO, C, N, and O. (The minor species NO, CO_2 , O^- , etc. were found to be negligible in the present calculation.)

FINAL ENTHALPY OF MIXTURE

The enthalpy per unit mass of the postcombustion mixture is given by

$$h_{II} = \sum_{i=1}^{9} n_i H_i = \frac{p}{\rho} + \sum_{i=1}^{9} n_i U_i$$
 (4)

where each $H_{\mathbf{i}}$ corresponds to the final temperature T of the mixture.

The mass per mole of postcombustion mixture is

$$M_{II} = \sum_{i=1}^{9} \frac{p_i}{p} M_i \tag{5}$$

Moreover, the number of moles of the ith species per unit mass of mixture is

$$n_{i} = \frac{1}{M_{TI}} \frac{p_{i}}{p} \tag{6}$$

Therefore, equations (4) and (6) yield

$$h_{II} = \frac{1}{M_{II}} \sum_{i=1}^{9} \frac{p_i}{p} H_i$$
 (7)

In accordance with the discussion centering around equation (3), the molar enthalpy H_5 of atomic cesium in the final mixture is

$$H_5 = \frac{5}{2} RT + U_5^{\text{(elec)}} - I$$
 (8)

where U_5 , the internal energy per mole of cesium due to electronic excitation, is given by (cf. appendix)

$$U_5^{\text{(elec)}} = RT^2 \frac{\partial}{\partial T} \log q_5^{\text{(elec)}}$$

$$=\frac{R}{k}\left(\frac{\epsilon_{1}e^{-\frac{\epsilon_{1}}{kT}}+2\epsilon_{2}e^{-\frac{\epsilon_{2}}{kT}}+2\epsilon_{3}e^{-\frac{\epsilon_{3}}{kT}}+3\epsilon_{4}e^{-\frac{\epsilon_{4}}{kT}}}{-\frac{\epsilon_{1}}{kT}+2e^{-\frac{\epsilon_{2}}{kT}}+2e^{-\frac{\epsilon_{3}}{kT}}+3e^{-\frac{\epsilon_{4}}{kT}}}\right)$$

The molar enthalpies of e and Cs are given by

$$H_8 = H_9 = \frac{5}{2} RT$$
 (9)

while the molar enthalpies H_1 , H_2 , H_3 , H_4 , H_6 , and H_7 of CN, N_2 , CO, C, N, and O, respectively, are given in reference 2.

In accordance with the assumption of adiabatic combustion at constant pressure, the mixture enthalpy per unit mass is the same before and after combustion; that is,

$$h_{T} = h_{TT} \tag{10}$$

where $h_{\rm I}$ and $h_{\rm II}$ are given by equations (2) and (7).

THERMODYNAMIC EQUILIBRIUM RELATIONS

The relative numbers of atoms present in the mixture follow from table I and are listed in table II.

TABLE II

Atom.	Relative number introduced
C N O Cs	$ \frac{(1 - \alpha)/(3 - 2\alpha)}{(1 - \alpha)/(3 - 2\alpha)} $ $ \frac{(1 - \alpha)/(3 - 2\alpha)}{\alpha/(3 - 2\alpha)} $

In following a standard procedure of physical chemistry, a fictitious "atomic" pressure $\,p_f\,$ is introduced, and table II is used to write the following equations relating the partial pressures $\,p_i\,$ of the principal species present in the postcombustion mixture. (These equations account for the conservation of all atoms of carbon, nitrogen, oxygen, and cesium, respectively.)

$$p_{1} + p_{3} + p_{4} = \frac{1 - \alpha}{3 - 2\alpha} p_{f}
 p_{1} + 2p_{2} + p_{6} = \frac{1 - \alpha}{3 - 2\alpha} p_{f}
 p_{3} + p_{7} = \frac{1 - \alpha}{3 - 2\alpha} p_{f}
 p_{5} + p_{8} = \frac{\alpha}{3 - 2\alpha} p_{f}$$
(11)

In the last equation, which represents the conservation of cesium, the partial pressure of Cs $^+$, namely p_9 , has been replaced by its equal, the partial pressure of e $^-$, that is, p_8 . Thus, the total pressure equation can be written

$$p_1 + p_2 + p_3 + p_4 + p_5 + p_6 + p_7 + 2p_8 = p$$
 (12)

The reactions to be considered (and their respective equilibrium constants K_i , i = 1, 2, 3, 5) are:

 $CN \rightleftharpoons C + N$ (dissociation of CN),

$$K_1 = \frac{p_4 p_6}{p_1} \tag{13}$$

 $N_2 \rightleftarrows 2N$ (dissociation of N_2),

$$K_2 = \frac{p_6^2}{p_2} \tag{14}$$

 $CO \stackrel{\rightarrow}{\leftarrow} C + O$ (dissociation of CO),

$$K_{3} = \frac{p_{4}p_{7}}{p_{3}} \tag{15}$$

 $Cs \stackrel{\rightarrow}{\leftarrow} Cs^+ + e^-$ (ionization of Cs),

$$K_5 = \frac{p_8 p_9}{p_5} = \frac{p_8^2}{p_5} \tag{16}$$

By using these expressions for the $\,\mathrm{K}_{\,\mathbf{i}}$, the partial pressures of the diatomic molecules and the cesium atoms can be eliminated from the four equations (11). Thus,

$$\frac{p_{4}p_{6}}{K_{1}} + \frac{p_{4}p_{7}}{K_{3}} + p_{4} = \frac{1-\alpha}{3-2\alpha} p_{f}$$

$$\frac{p_{4}p_{6}}{K_{1}} + 2 \frac{p_{6}^{2}}{K_{2}} + p_{6} = \frac{1-\alpha}{3-2\alpha} p_{f}$$

$$\frac{p_{4}p_{7}}{K_{3}} + p_{7} = \frac{1-\alpha}{3-2\alpha} p_{f}$$

$$\frac{p_{8}p_{7}}{K_{5}} + p_{8} = \frac{\alpha}{3-2\alpha} p_{f}$$
(17)

Values for the equilibrium constants K_1 , K_2 , and K_3 are given in reference 2. Note, however, that the definition of K_2 in equation (14) differs from the equilibrium constant defined for the reaction $N_2 \rightleftarrows 2N$ in reference 2. The equilibrium constant K_5 is given by

$$K_{5} = \frac{2}{q_{5}} \left(\frac{2\pi}{h^{2}} \frac{m_{8}m_{9}}{m_{5}} \right)^{3/2} (kT)^{5/2} e^{-\frac{I'}{kT}}$$
(18)

Equation (18) is derived in the appendix.

SOLUTION OF THE SIMULTANEOUS EQUATIONS

Equations (10), and (12) to (17) constitute ten equations in the ten unknowns p_1 , p_2 , ... p_8 , p_f , and T. The total pressure p_8 and the initial cesium mole fraction α enter as parameters. (The fictitious pressure p_f , which could have been eliminated immediately, affords the convenience of a Lagrangian undetermined multiplier.)

A practicable procedure for solving the simultaneous equations is the following:

- (1) Select the total pressure p and the initial cesium mole fraction α .
- (2) Estimate the partial pressures p_i (i = 4, 6, 7), fictitious pressure p_f , and temperature T.
- (3) Solve the first three of equations (17) for p_4 , p_6 , and p_7 by an extension of Newton's method of successive approximation. (In the event that the convergence is slow because the iterative character is oscillatory, it may be advantageous to use corrections less than the full Newtonian ones.)
- (4) Compute the corresponding values of the remaining partial pressures from the fourth of equations (17) and equations (13) to (16).
- (5) Equation (12) generally will not be satisfied. Repeat steps 3 and 4 with a new value of $p_{\rm f}$ until equation (12) is satisfied.

(6) Equation (10) generally will not be satisfied. Repeat steps 3, 4, and 5 with a new value of T until equation (10) is satisfied.

The computing time is much shortened by making careful estimates of the trial values. Extrapolations become possible as the computation proceeds.

RESULTS OF THE CALCULATION

The calculation has been carried through for four values of the total pressure and six values of the initial mole fraction of cesium, namely, p = 1, 0.2, 0.05, 0.01 atm and α = 0, 0.02, 0.04, 0.06, 0.08, and 0.1. The results are presented in table III and figures 1 to 3. In addition to the postcombustion temperature and partial pressures of the principal species the electron concentration

$$N_e = \frac{p_{\theta}}{kT}$$

has been included in the table in cgs units (electrons cm⁻³). (To convert p_8 from atm to dynes/cm², multiply by 1.01325 \times 10⁶.)

Figure 1 shows that the addition of cesium to the extent of 0.1 mole fraction results in a lowering of the flame temperature by approximately 300° K at each of the four pressures considered. The greatest electron concentrations in the cases considered occur when p=1 atm, but somewhat greater efficiency of ionization of the cesium is obtained at the lower pressures. To a first approximation, the principal species present in the postcombustion mixture (aside from Cs, Cs⁺, and e⁻) are CO and N₂, followed in decreasing prevalence by N, O, and so forth. As the total pressure is lowered from 1 atm to 0.01 atm the temperature drops by approximately 500° K for any fixed value of α in the range considered.

The maximum electron concentration obtainable at p=1 atm must occur for α greater than 0.1 and probably is of the order of 5×10^{16} electrons cm⁻³ (cf. fig. 2), but the computations have not been extended to verify this.

Figure 3 shows that for small values of α and for low pressures, the electron mole fraction is independent of the pressure. In such cases, the ionization is practically complete. (See table III.)

Langley Research Center,
National Aeronautics and Space Administration,
Langley Field, Va., February 10, 1960.

APPENDIX

DERIVATION OF EXPRESSION FOR K5 FROM SAHA'S EQUATION

In a given volume V, let

 v_5 = number of Cs atoms

 v_8 = number of free electrons

 v_{Q} = number of Cs⁺ ions

I' = ionization energy per Cs atom

Then (cf. ref. 3),

$$\frac{v_8 v_9}{v_5} = \frac{q_8 q_9}{q_5} e^{-\frac{I'}{kT}}$$

where

$$q_8 = 2 \frac{(2\pi m_8 kT)^{3/2}}{h^3} v$$

$$q_9 = \frac{(2\pi m_9 kT)^{3/2}}{h^3} V$$

$$q_{5} = q_{5}^{\text{(elec)}} \frac{(2\pi m_{5} kT)^{3/2}}{h^{3}} V$$

$$q_{5}^{\text{(elec)}} = 2 + 2e^{-\frac{\epsilon_{1}}{kT}} + 4e^{-\frac{\epsilon_{2}}{kT}} + 4e^{-\frac{\epsilon_{3}}{kT}} + 6e^{-\frac{\epsilon_{4}}{kT}}$$

But,

$$p_i V = v_i kT$$
 (i = 5, 8, 9)

Hence,

$$\frac{v_8 v_9}{v_5} = \frac{p_8 p_9}{p_5} \frac{v}{kT} = K_5 \frac{v}{kT}$$

and

$$K_{5} = \frac{\left(q_{8}/V\right)\left(q_{9}/V\right)}{q_{5}/V} \text{ kTe}^{-\frac{I'}{kT}}$$

$$= \frac{2}{q_{5}^{(elec)} \left(\frac{2\pi}{h^{2}} \frac{m_{8}m_{9}}{m_{5}}\right)^{3/2} (kT)^{5/2} e^{-\frac{I'}{kT}}$$

which is the desired equation (18).

The preceding derivation accounts for the excited electronic states of the cesium atom but neglects those of the cesium ion since the latter are practically inaccessible in the temperature range of interest here.

REFERENCES

- 1. Conway, J. B., Smith, W. F. R., Liddell, W. J., and Grosse, A. V.: The Production of a Flame Temperature of 5000° K. Jour. American Chem. Soc. (Communications to Ed.), vol. 77, no. 7, Apr. 5, 1955, pp. 2026-2027.
- 2. Gordon, John S.: Thermodynamics of High-Temperature Gas Mixtures, and Application to Combustion Problems. WADC Tech. Rep. 57-33, ASTIA Doc. No. 110735, U.S. Air Force, Jan. 1957.
- 3. Fowler, R. H., and Guggenheim, E. A.: Statistical Thermodynamics. The Macmillan Co., 1939.

TABLE III.- TEMPERATURE AND COMPOSITION OF CYANOGEN-OXYGEN FLAME SEEDED WITH CESTUM

L-605

		V9	10		+
Ne,		1016	1015		1014
		1.0.0.4. 5.4.0.4. X	×	00890	×
		10004	6.98 6.99 7.11 7.74	- คี่ ผู้ ผู้ ค้า	04.08 OU
Cs+		10 21 24 27	.0024 .0043 .0058 .0059	0006 0012 0017 0021 0024	.00012 .00025 .00037 .00047
Ö		0.010 .017 .021 .021	, , , , , , , , , , , , , , , , , , ,	88888	ŏŏŏŏŏ
		9 7 7 7 7	0024 0043 0058 0069	.0006 .0012 .0017 .0021	00012 00025 00037 00047 00067
' o		0.010 .017 .021 .024 .027	88888	88888	88888
		94450	11 11 06 05	05 02 02 01	00012 000010 00008 00006 00005
0		0.006 .004 .005 .005 .007	.0014 .0014 .0011 .0006 .0006	.0005 .0007 .0002 .0002 .0002	888888
	ដ	5 6 5 5	25 25 25 16 13	113 008 05 05 05	00034 00028 00027 00017 00014 00010
N	, atm	0.013 .010 .008 .007 .006	.00400. .0032 .0025 .0020 .0016	.0013 .0001 .0008 .0007 .0005	888888
	Partial pressure,	20800	22522	84558	000 000 000 000 000 000 000 000 000 00
Çs	pre	0.003 .010 .018 .029 .040	.0002 .0010 .0021 .0036	.000 .000 .000 .000 .000 .000	.00001 .00001 .00002 .00004
	tial		4 H& V V A	ちょうとらは	250 045 03 03 03
ũ	Par	.005 .004 .003 .002 .002	.0011 .0011 .0008 .0006 .0007	.0005 .0004 .0003 .0002	.00012 .00009 .00006 .00006
		0	0 > + 0 0 >	0-48 B+B	18815 <u>*</u>
8		652 640 628 628 607 607	.1290 .1277 .1244 .1222 .1202	.0321 .0314 .0308 .0303 .0297	.00622 .00622 .00601 .00600 .00588
		0			0 4 5 4 0 10
N2		. 322 . 317 . 311 . 307 . 301 . 301	0622 0628 0613 0596 0587	.0156 .0157 .0157 .0159 .0141	.00506 .00501 .00297 .00294 .00285
		0			
CN		.002 .001 .001 .001	0003 0002 0002 0002 0002	0000 0000 0000 0000	00001 00001 00001 00000
		0			
ð	4	< 103			
1	۲,	4.85 × 4.52 × 4.52	4.70 4.54 4.47 4.47	4.56 4.41 4.41 4.33 4.26	4.38 4.32 4.25 4.19 4.19
	ಕ	0.00 40.00 60.00 10.00 10.00	× 5,5,5,5,	95,000	20000
p,	atm	1.00	8.	.05	<u>></u>
				•	•

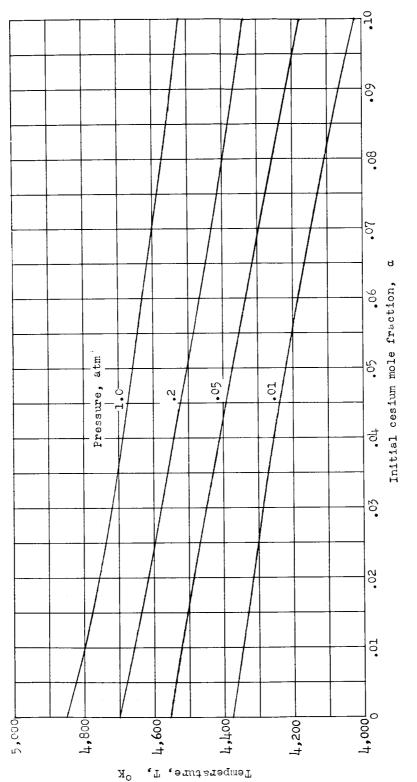
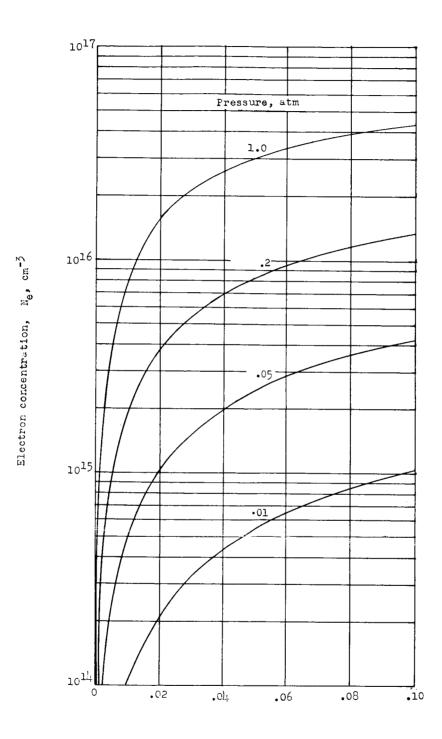


Figure 1.- Variation of temperature with initial cesium mole fraction for various values of total pressure.

L-605



Initial cesium mole fraction, α

Figure 2.- Variation of electron concentration with initial cesium mole fraction for various values of total pressure.

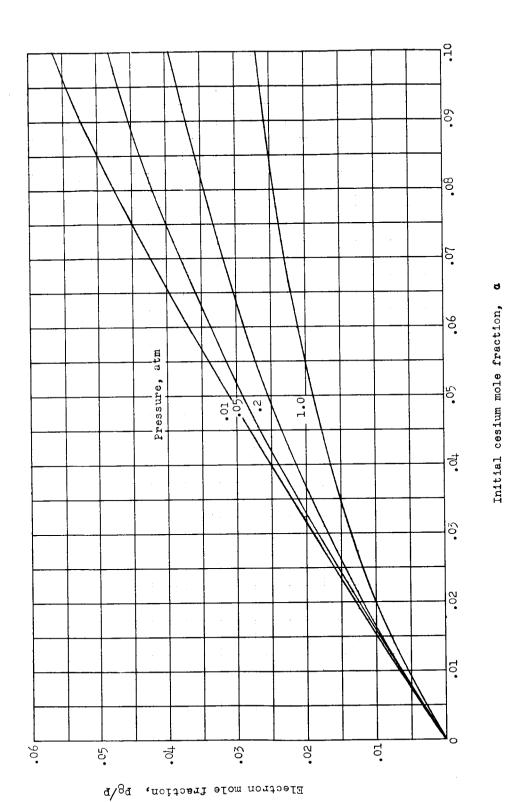


Figure 3.- Variation of electron mole fraction with initial cesium mole fraction for various values of total pressure.